

Ligand Field Analyses of Tris(biuret)chromium(III) Chloride and Hexaureachromium(III) Bromide

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Ligand field analyses for tris(biuret)chromium(III) chloride and hexaureachromium(III) bromide were performed and compared to understand the ligand field properties of both ligands. The optimized $e_{\sigma 0}$ and $e_{\pi 0}$ values indicate that coordinated oxygen atom in biuret ligand is a moderate σ - and strong π -donor, while that in urea ligand is a weak σ - and moderate π -donor. The electronic structures of those two complexes are quite different and they were well accounted by inclusion of an anisotropic π bonding.

Introduction

The spectroscopic properties of $[\text{Cr}(\text{urea})_6]^{3+}$ complex have been extensively studied.¹⁻⁸ The close proximity of doublet (2E_g) and quartet (${}^4T_{2g}$) states of this complex give rise to many unusual photophysical properties and these features have not been satisfactorily interpreted. Though the knowledge of ligand field properties and electronic structure is essential to understand this system, its ligand field analysis has been largely neglected, however.

It has been reported that its chelating system, $[\text{Cr}(\text{biuret})_3]^{3+}$ exhibit quite different photophysical properties as well as ligand field properties from those of $[\text{Cr}(\text{urea})_6]^{3+}$ complex.^{9,10} For one instance, $[\text{Cr}(\text{biuret})_3]^{3+}$ shows larger ligand field strength and resulting different energy gap between 2E_g and ${}^4T_{2g}$ states of the two complexes exhibit discriminatory emission behavior; only phosphorescence is observed in the biuret complex, while phosphorescence-fluorescence dual emissions were observed in the urea complex. Since the symmetry of these complexes is trigonal (D_3), symmetry lowering of biuret complex upon chelation cannot be a major factor for this discrepancy.⁹ Moreover, doublet line splittings are more sensitive in urea complex to the subtle geometry changes induced by counter ions. Thus, it is appropriate to analyze and compare the ligand field properties of each ligand to understand the electronic structure of these CrO_6 coordination system.

In this paper, we examined the spectroscopic properties of $[\text{Cr}(\text{biuret})_3]^{3+}$ and performed the ligand field analyses of $[\text{Cr}(\text{biuret})_3]^{3+}$ and $[\text{Cr}(\text{urea})_6]^{3+}$ to elucidate electronic structures and ligand field properties of the ligands. In the framework of the angular overlap model (AOM), the suitability of anisotropic π -bonding was examined in reproducing the unusual doublet splittings in these complexes.

Experimental Section

Tris(biuret)chromium(III) chloride was prepared by a standard method.¹⁰ The needle-shaped, deep green crystals

were obtained from successive recrystallization in 0.1M HCl solution in the dark. Luminescence and excitation spectra were measured with a Continuum Nd : YAG laser-pumped dye laser (ND60) and BMI OPO laser as a light source. The emission was analyzed with a CVI 0.5m monochromator (DK480), and was detected with a Hamamatsu R943-02 cooled photomultiplier tube and an SRS boxcar chassis. Microcrystalline samples were mounted on the cold head of a Janis CCS-600 closed cycle He gas cryostat. IR and Raman spectra were measured with a Perkin-Elmer System 2000 FT-IR spectrometer on samples dispersed in Nujol mulls on a polyethylene film (far-IR) or in KBr pellets (mid-IR). Room temperature absorption spectra were recorded with a Shimadzu UV3100 spectrophotometer.

Results and Discussion

Vibrational spectra. The 13K luminescence spectrum of $[\text{Cr}(\text{biuret})_3]\text{Cl}_3$ is shown in Figure 1. Peaks are rather broad and vibrational side bands are weak, compared to other Cr^{3+} complexes.^{5,11} No trace of impurities was detected in temperature and exciting wavelength dependent spectroscopy. The strong peak at 14245 cm^{-1} is coincident with a strong, lowest energy peak in excitation spectrum and

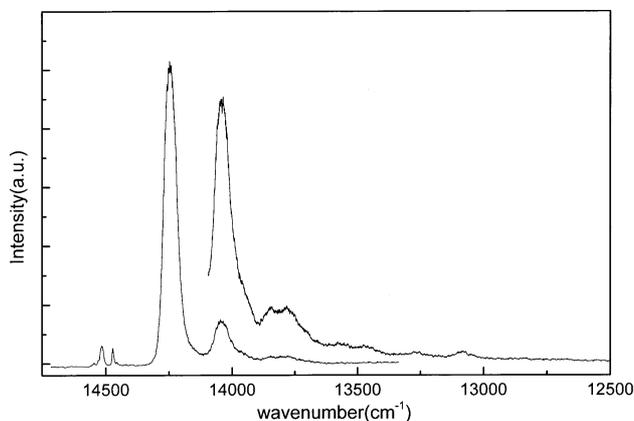


Figure 1. Luminescence spectrum of $[\text{Cr}(\text{biuret})_3]\text{Cl}_3$ at 13K.

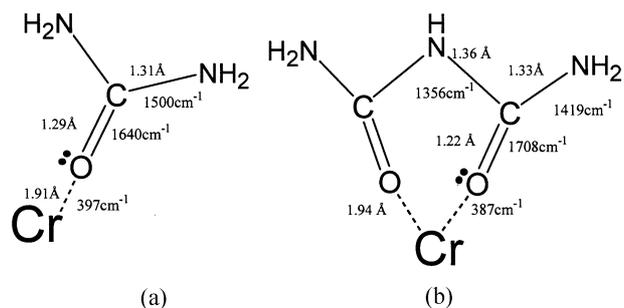
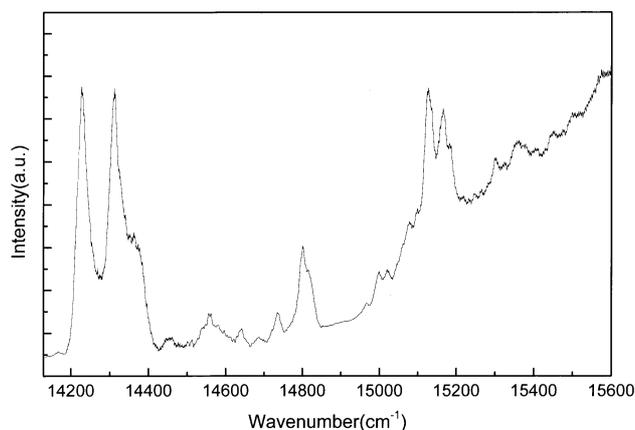
Table 1. Vibronic intervals in the 13K luminescence spectrum of $[\text{Cr}(\text{biuret})_3]\text{Cl}_3$

| Vibrational Interval(cm^{-1}) | IR | Raman | Assignment | $[\text{Cr}(\text{urea})_6]^{3+}$ Ref. 5 | Biuret Ligand |
|--|-----|-------|--------------------|--|---------------|
| -297 | | | | | |
| -269 | | | | | |
| -225 | | | Hot bands | | |
| -208 | | | | | |
| -93 | | | | | |
| -63 | | | | | |
| 0 | | | 2E_g origin | | |
| 25 | | | Lattice | | |
| 67 | | | Lattice | | |
| 106 | 99 | 102 | Rocking or torsion | | |
| | 116 | 114 | | | |
| 210 | 206 | 206 | (O-Cr-O) | 198 | |
| | | | | 207 | |
| 229 | 222 | 226 | Sym (O-Cr-O) | 226 | |
| | | | | 243 | |
| 272 | 264 | 276 | (O-Cr-O) | | |
| | 282 | | | | |
| 298 | 298 | | Asym (Cr-O) | 294 | |
| | 325 | 329 | (Cr-O) | 327 | |
| 387 | 377 | 383 | Sym (Cr-O) | 397 | |
| | 434 | 434 | Ligand | | 439 |
| 463 | | | 229 2 (458) | | |
| | 467 | 465 | Ligand | | 460 |
| | 505 | 505 | Ligand | | 500 |
| 685 | | | 387 + 298(685) | | |
| 769 | | | 387 2 (774) | | |
| 978 | | | | | |

assigned to the lower energy component of 2E_g origin. The luminescence spectrum of perchlorate salt is very similar to that of chloride salt and the 2E_g origin is found at 14227 cm^{-1} .

Vibronic intervals in the luminescence spectrum along with IR and Raman frequencies are listed in Table 1. Though the approximate symmetry of $[\text{Cr}(\text{biuret})_3]\text{Cl}_3$ is D_3 , the actual symmetry is much lower and the symmetry of vibrational modes was not well represented in the two complementary techniques, Raman and IR. The symmetry of vibrational modes was estimated by comparing the relative intensities of the corresponding bands in both spectra. The bands below 250 cm^{-1} correspond to lattice and ligand deformation modes with considerable O-Cr-O bending character. A weak peak at 387 cm^{-1} appears in IR and Raman with considerable intensity and assigned to the symmetric Cr-O stretching mode. A weak peak at 298 cm^{-1} was absent in Raman spectrum and assigned to the asymmetric Cr-O stretching mode. Above 400 cm^{-1} , the frequency of these bands can be found in the free ligand IR spectrum and they are assignable to internal modes of ligands. As seen in Table 1, the frequencies of Cr-O vibration and ligand deformation modes are similar between hexaurea and biuret complexes.

The mid-IR spectrum shows minor frequency changes in most of the biuret ligand vibrations upon complexation with

**Figure 2.** Comparison of IR stretching frequencies and bond lengths (a) $\text{Cr}(\text{urea})_6^{3+}$ (b) $\text{Cr}(\text{biuret})_3^{3+}$.**Figure 3.** Excitation spectrum of $\text{Cr}(\text{biuret})_3\text{Cl}_3$ at 13K.

Cr metal, in contrast to the urea complex. The C-N stretching frequencies remain unchanged and the carbonyl stretching frequencies decrease by a few wavenumbers upon complexation in the biuret complex while considerable decrease of $\nu(\text{C}=\text{O})$ and increase of $\nu(\text{C}-\text{N})$ frequencies are observed in the urea complex.¹²

The vibrational frequencies and bond lengths of urea and biuret complexes are compared in Figure 2. The x-ray crystal structure of urea suggested the resonance structure of 30% double bond character for the carbon to nitrogen bonds and 40% for the remaining carbon to oxygen bond.¹³ Biuret ligand lacks this type of resonance and the double bond character remains mostly in the carbonyl moiety of the ligand.

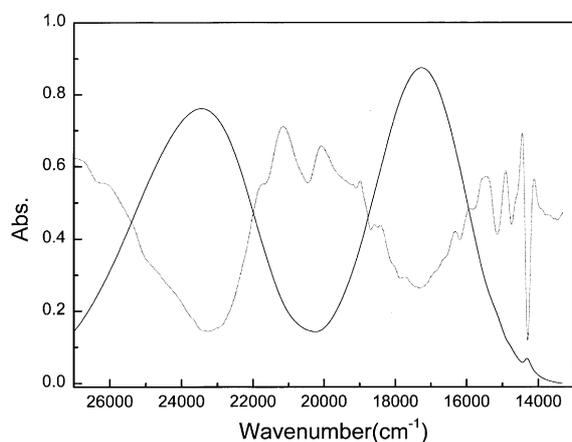
Electronic spectra. The 12K excitation spectrum of $[\text{Cr}(\text{biuret})_3]\text{Cl}_3$ is shown in Figure 3 and the peak position and their assignments are tabulated in Table 2. The vibronic lines are also quite attenuated, as seen in the luminescence spectrum. The electronic origins are expected to be quite strong since the effective symmetry of the chromophore is very low. A strong peak at 14245 cm^{-1} coincides with the peak in the luminescence spectrum and is assigned to lower energy component of 2E_g origin. The other component of 2E_g origin is found at 14329 cm^{-1} . The 84 cm^{-1} of the 2E_g splitting is comparable to that of $\text{Cr}(\text{urea})_6^{3+}$ complex. Three components of $^4A_{2g} \rightarrow ^2T_{1g}$ transition can be found with strong intensities at $573, 900, 937 \text{ cm}^{-1}$ from the lowest electronic line. As seen in Figure 3, the overlap of the $^2T_{1g}$ transition region with the broad $^4T_{2g}$ band and the consequent

Table 2. Peak positions in the 13K excitation spectrum of [Cr(biuret)₃]Cl₃. (all data in cm⁻¹)

| $\bar{\nu}$ - 14245 | Assignment | $\bar{\nu}$ - 14245 | Assignment |
|---------------------|-------------------------------------|---------------------|--------------------------------------|
| 0 s | ² E _g | 793 w | ² T _{1g} + 229 |
| 29 w,s h | ² E _g + 25 | 850 vw | ² T _{1g} + 272 |
| 66 vw, sh | ² E _g + 67 | 871 vw | ² T _{1g} + 298 |
| 84 s | ² E _g ' | 900 m | ² T _{1g} ' |
| 100 vw, sh | ² E _g + 106 | 937 m | ² T _{1g} " |
| 136 w, sh | | 958 w, sh | ² T _{1g} + 387 |
| 228 vw, sh | ² E _g + 229 | 1008 vw | ² T _{1g} ' + 106 |
| 274 vw | ² E _g + 272 | 1018 vw | |
| 284 vw | ² E _g ' + 210 | 1037 vw | ² T _{1g} " + 106 |
| 314 vw, sh | ² E _g ' + 229 | 1074 w | |
| 349 w | ² E _g ' + 272 | 1097 vw | ² T _{1g} ' + 210 |
| 382 vw, sh | ² E _g + 387 | 1143 w | ² T _{1g} " + 210 |
| 412 vw | | 1151 vw | ² T _{1g} " + 229 |
| 469 vw | ² E _g ' + 387 | 1181 vw | ² T _{1g} ' + 272 |
| 509 vw | | 1224 w | ² T _{1g} " + 272 |
| 573 m | ² T _{1g} | 1275 vw | ² T _{1g} ' + 387 |
| 738 vw | | 1337 vw | ² T _{1g} " + 387 |
| 770 w | ² E _g ' + 685 | 1372 vw | ² T _{1g} ' + 463 |

intensity borrowing by the doublet lines are apparent. However, none of the well resolved peaks appear to belong to the ⁴T_{2g} band. These doublet lines are also observed in room temperature solution absorption and its second derivative spectrum as seen in Figure 4. The components of ²E_g and ²T_{1g} transitions are located at 14305, 14747 and 15156 cm⁻¹ and these support above assignment. The relatively strong band at 20833 cm⁻¹ is assigned to the first components of the ²T_{2g} transitions, but the spectrum is not resolved well enough to assign the other components with any certainty.

The broad spin allowed ⁴T_{2g} and ⁴T_{1g} bands are observed at 17301 and 23725 cm⁻¹. Split quartet components in D₃ symmetry could not be found even in the second derivative spectrum, while they are found in polarized absorption spectrum of urea complex.¹⁴ It can be explained by the possibility that they are closely located with each other, or the intensity of some components (⁴E from ⁴T_{2g}, ⁴A₂ from ⁴T_{1g}, for example) is so weak that they are buried in other com-

**Figure 4.** Absorption spectrum (solid line) and second derivative spectrum (dotted line) of Cr(biuret)₃³⁺ in 0.1M HCl solution at room temperature.

ponents. It is interesting to note that the ligand field strength and effective ligand field symmetry increase with chelation. The quartet bands of biuret complex are shifted to higher energy by about 1200 cm⁻¹ relative to those of the hexaureachromium(III) complex.

Ligand field analysis. Ligand field analyses for [Cr(biuret)₃]³⁺ complex as well as hexaurea complex were performed to understand the ligand field properties of the two ligands. The general method of ligand field calculations in the framework of angular overlap model has been described elsewhere.^{15,16} The ligand field potential matrix was generated from the six coordinated oxygen atoms, by use of the room temperature X-ray single crystal structure.^{9,17-19}

The six parameters varied during the optimization were the AOM ligand field parameters $e_{\sigma O}$ and $e_{\pi O}$, plus interelectronic repulsion parameters B , C and α_T (the Trees correction parameters), and the spin-orbit coupling parameter, ζ . The π interaction of the oxygen with the metal ion was considered to be anisotropic, confined to the direction perpendicular to the plane defined by the Cr-O-C bonds. The experimental energies used in the fitting process, along with their assignments, are given in Table 3. By variation of six parameters just described, these energies were fit by means of the Powell parallel subspace optimization procedure.²⁰ The function minimized was

$$F = \sum Q^2 + 10 \sum T^2 + 100 \sum D^2 + 1000 \sum S^2$$

where each term represents a difference between experimental and calculated transition energies or splittings: D , the five lowest doublet energies; T , the averaged ²T₂ peak position; Q , the two quartet energies and S , the splittings between the doublet energies. The weighting factors in this function are in approximate proportion to the inverse square of the corresponding experimental uncertainty. Reasonable boundary conditions based on the results from other Cr(III) complexes were applied to each parameter, but none of the parameters in the best-fit parameter set approached its boundary value. The optimization was repeated several times with different sets of starting values to confirm that the same global mini-

Table 3. Experimental and calculated transition energies for Cr(biuret)₃³⁺ and Cr(urea)₆³⁺ complex. (all data in cm⁻¹)

| | Cr(biuret) ₃ ³⁺ | | Cr(urea) ₆ ³⁺ | |
|------------------------------------|---------------------------------------|--------------------|--|--------------------|
| | Obs. | Calc. ^a | Obs. ^{c,d} | Calc. ^b |
| ² E _g | 14245 | 14229 | ² E _g 14184 | 14182 |
| | 14329 | 14348 | | 14265 |
| | 14818 | 14874 | ² T _{1g} 14361 | 14365 |
| ² T _{1g} | 15145 | 15092 | | 14963 |
| | 15182 | 15152 | | 15129 |
| | 20833 | 21603 | ² T _{2g} 20762 | 20749 |
| ⁴ T _{2g} (avg) | 17301 | 17287 | ⁴ T _{2g} ⁴ E 16130 ^d | 16118 |
| ⁴ T _{1g} (avg) | 23725 | 23750 | ⁴ A ₁ 17000 | 17217 |
| | | | ⁴ T _{1g} ⁴ E 22570 | 22306 |
| | | | ⁴ A ₂ 23400 | 23133 |

^aligand field parameters: $e_{\sigma O} = 7055$, $e_{\pi O} = 1876$, $B = 689$, $C = 2910$, $T = 106$, $\zeta = 24$. ^b $e_{\sigma O} = 6200$, $e_{\pi O} = 1406$, $B = 676$, $C = 3121$, $T = 156$, $\zeta = 194$. ^cspectrum of bromide salt, Ref. 5. ^dRef. 14.

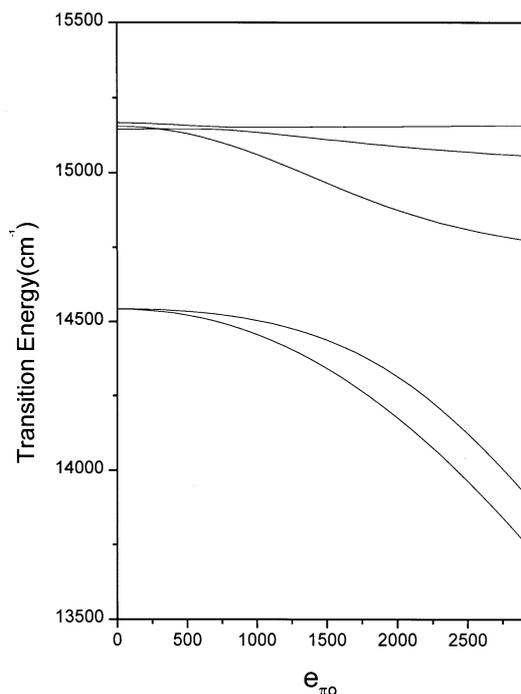


Figure 5. Calculated variation of the transition energies to the 2E_g and ${}^2T_{1g}$ state for $\text{Cr}(\text{biuret})_3^{3+}$ as a function of $e_{\pi o}$.

mum was found.

The results of the optimization and the parameter set used to generate the best-fit energies are listed in Table 3. The optimized $e_{\sigma o}$ and $e_{\pi o}$ values indicate that coordinated oxygen atom in biuret ligand is a medium σ - and strong π -donor, while that in urea ligand is a weak σ - and medium π -donor. It is in accordance with the results of x-ray crystal structure and vibrational analysis, as mentioned in previous section. The electron density is lower in carbonyl moiety of the urea ligand than that of a biuret since urea involves considerable resonance hybrid. The σ -bonding contribution is judged to be more important in determining the ligand field strength of ligands.

Table 3 presents Flint's data for urea complex, which are noteworthy for the large overall doublet splittings, especially in ${}^2T_{1g}$ region, and the close proximity of 2E_g and ${}^2T_{1g}$ states. This is quite surprising since we expect symmetry is lowered upon chelation and larger overall splitting in biuret complex is expected. Indeed, molecular symmetry of biuret complex is lower than that of hexaurea complex.⁹ However, both urea and biuret are non-linear ligand and the orientation of their π orbital may affect the metal d orbital energies significantly. By close inspection of x-ray crystal data, the biuret complex shows the π orbital orientation angle ψ to be 40° , which implies pseudoisotropic π -bonding, while the urea complex shows anisotropic π -bonding in which ψ varies from 58° to 72° . The isotropic condition results when $e_{\pi s}$ is equal to $e_{\pi c}$ at $\psi = 45^\circ$, since the overlap of the ligand π orbital with metal d_{xz} and d_{yz} orbitals is then equal.¹³

This geometric difference leads to a quite different behavior of the doublet transitions. Figures 5 and 6 show the calculated behavior of the doublet transition energies as a

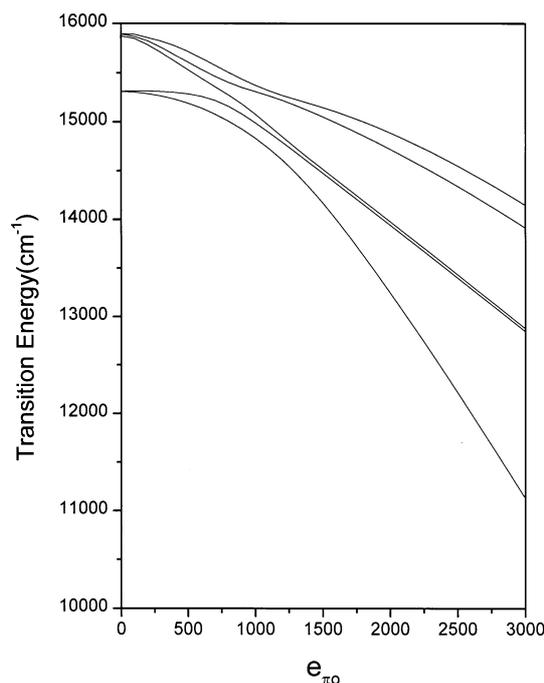


Figure 6. Calculated variation of the transition energies to the 2E_g and ${}^2T_{1g}$ states for $\text{Cr}(\text{urea})_6^{3+}$ as a function of $e_{\pi o}$.

function of the extent of anisotropic π -bonding of biuret and urea complexes, respectively. The $e_{\sigma o}$ was adjusted to keep the quartet energies approximately constant. The doublet splittings are quite insensitive to the value of e_{π} in the pseudoisotropic biuret system while doublet splittings become larger and show the considerable mixing between the ${}^2T_{1g}$ and 2E_g states as the $e_{\pi o}$ increase in the anisotropic urea system. At $e_{\pi o}$ of around 1500 cm^{-1} , the overall doublet splitting pattern of 1-2-2 is reproduced and fit the experimental data of Flints quite well. The second component of the ${}^2T_{1g}$ transition located at 14541 cm^{-1} deviate only from calculated energies considerably, however. This may be a vibronic line and the calculated energies suggest that the second component of the ${}^2T_{1g}$ transition line should be found near 15000 cm^{-1} . The inclusion of an anisotropic bonding markedly improve the fitting and its effect on the doublet line splittings are substantial for these complexes.

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References

- Porter, G. B.; Schläfer, H. L. *Ber. Bunsenges. Physik. Chem.* **1964**, *68*, 316.
- Gausmann, H.; Schläfer, H. L. *J. Chem. Phys.* **1968**, *48*(9), 4056.
- Dingle, R. *J. Chem. Phys.* **1969**, *50*, 1592.
- Yersin, H.; Otto, H.; Gliemann, G. *Theor. Chim. Acta.* **1974**, *33*, 63.
- Flint, C. D.; Palacio, D. J. D. *J.C.S. Faraday II* **1979**, *75*, 1159.

6. Flint, C. D.; Palacio, D. J. D. *J.C.S. Faraday II* **1980**, 76, 82.
 7. Yersin, H.; Huber, P.; Gietl, G.; Trümbach, D. *Chem. Phys. Lett* **1992**, 199(1), 1.
 8. Schmidtke, H. H.; Schoepe, K. E.; Degen, J. *Ber. Bunsenges. Physik. Chem.* **1995**, 99, 759.
 9. Park, S. J.; Lee, K.W.; Park, Y. J. To be published.
 10. Chatterjee, K. K.; Porter, G. B. *Inorg. Chem.* **1966**, 5, 5, 860.
 11. Mortensen, O. S. *J. Chem. Phys.* **1967**, 47(10), 4215.
 12. Penland, R. B.; Mizushima, S.; Curran, C.; Quagliano, J. V. *J. Am. Chem. Soc.* **1957**, 79, 1575.
 13. Vaughan, P.; Donohue, J. *Acta Cryst.* **1952**, 5, 530.
 14. Schönherr, T. *Topics in Curr. Chem.* **1997**, 191, 87.
 15. Hoggard, P. E. *Coord. Chem. Rev.* **1986**, 70, 85.
 16. Lee, K. W.; Eom, K. I.; Park, S. J. *Inorg. Chim. Acta* **1997**, 254, 131.
 17. Davis, P. H.; Wood, J. S. *Inorg. Chem.* **1970**, 9(5), 1111.
 18. Okaya, Y.; Pepinsky, R. *et al. Acta. Cryst.* **1957**, 10, 798.
 19. Figgis, B. N.; Wadley, L. G. B.; Graham, J. *Acta Cryst.* **1972**, B28, 187.
 20. Powell, M. J. D. *Comput. J.* **1964**, 7, 155.
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