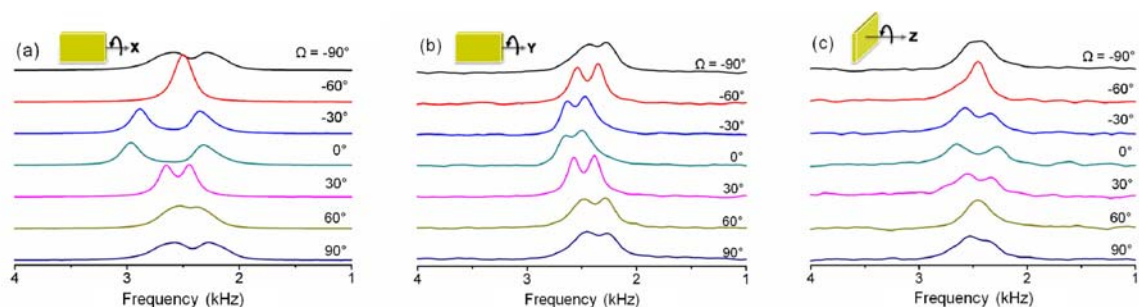
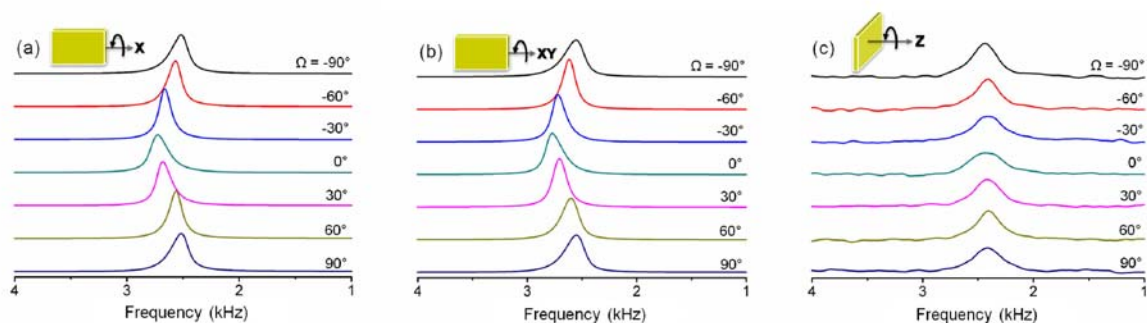


## Supporting Information

Hot-Pressing Effects on Polymer Electrolyte Membrane Investigated by  $^2\text{H}$  NMR SpectroscopySang Man Lee<sup>a</sup> and Oc Hee Han<sup>†,‡,\*</sup>*Daegu Center, Korea Basic Science Institute, Daegu 702-701, Korea. \*E-mail: ohhan@kbsi.re.kr**<sup>†</sup>Graduate School of Analytical Science & Technology, Chungnam National University, Daejeon 305-764, Korea**<sup>‡</sup>Department of Chemistry, Kyungpook National University, Daegu 702-701, Korea**Received December 22, 2012, Accepted December 27, 2012*

**Figure S1.** Representative spectra of the stretched Nafion 117 membranes obtained at 3 different rotation axes as designated in the spectra. The  $^2\text{H}_2\text{O}$  contents were 13, 14, and 13 wt% for the membranes rotated along the (a) X, (b) Y, and (c) Z axes, respectively.



**Figure S2.** Representative spectra of the Nafion 117 membranes without any stress applied obtained at 3 different rotation axes as designated in the spectra. The  $^2\text{H}_2\text{O}$  contents were (a) 23 wt% for the sample rotated along the X axis, (b) 29 wt% for the sample rotated along the XY axis, and (c) 3 wt% for the sample rotated along the Z axis.

**[Relationships among  $\theta$ ,  $\Omega$ ,  $\alpha$ , and  $\beta$ ]**

The average O-<sup>2</sup>H director and H<sub>0</sub> can be expressed as vectors **A** and **B**, respectively, in the XYZ frame, adopting the method employed in reference S1. Using the vector product  $\mathbf{A} \cdot \mathbf{B} = |\mathbf{A}||\mathbf{B}|\cos\theta$ , the relationship between angles can be obtained if **A** and **B** were considered unit vectors.

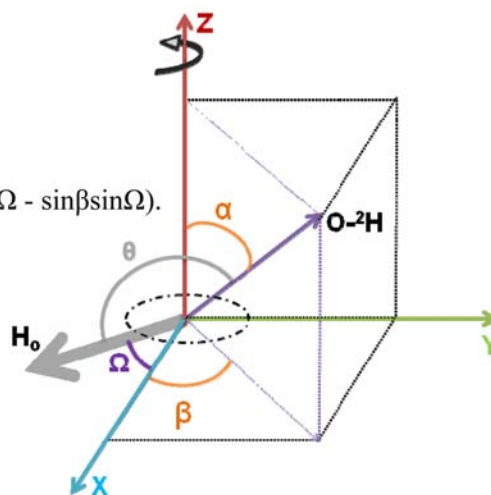
**For the Z rotation**

$$\mathbf{A} = (\sin\alpha\cos\beta, \sin\alpha\sin\beta, \cos\alpha)$$

$$\mathbf{B} = (\cos\Omega, -\sin\Omega, 0),$$

resulting in

$$\cos\theta = \sin\alpha\cos\beta\cos\Omega - \sin\alpha\sin\beta\sin\Omega = \sin\alpha(\cos\beta\cos\Omega - \sin\beta\sin\Omega).$$

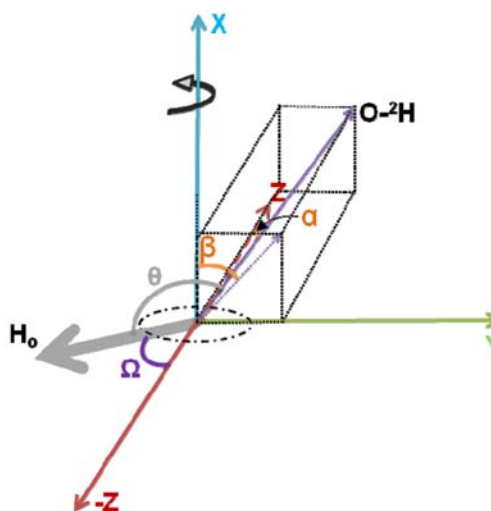
**For the X rotation**

$$\mathbf{A} = (\sin\alpha\cos\beta, \sin\alpha\sin\beta, \cos\alpha)$$

$$\mathbf{B} = (0, -\sin\Omega, -\cos\Omega),$$

resulting in

$$\cos\theta = -\sin\alpha\sin\beta\sin\Omega - \cos\alpha\cos\Omega.$$



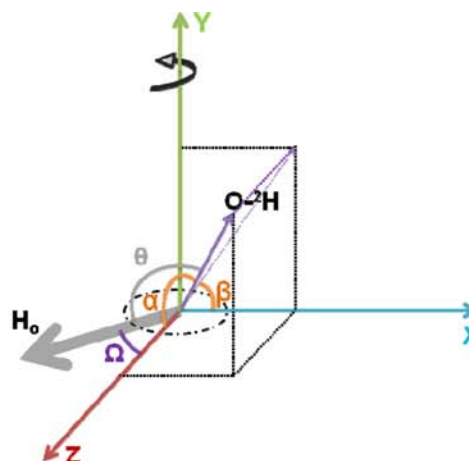
**For the Y rotation**

$$\mathbf{A} = (\sin\alpha\cos\beta, \sin\alpha\sin\beta, \cos\alpha)$$

$$\mathbf{B} = (-\sin\Omega, 0, \cos\Omega),$$

resulting in

$$\cos\theta = -\sin\Omega\sin\alpha\cos\beta + \cos\Omega\cos\alpha.$$

**For the XY rotation**

$$\mathbf{A} = (\sin\alpha\cos\beta, \sin\alpha\sin\beta, \cos\alpha)$$

$$\mathbf{B} = (\sin\Omega\cos(45^\circ), -\sin\Omega\sin(45^\circ), \cos\Omega) = (0.707\sin\Omega, -0.707\sin\Omega, \cos\Omega),$$

resulting in

$$\cos\theta = 0.707\sin\alpha\cos\beta\sin\Omega - 0.707\sin\alpha\sin\beta\sin\Omega + \cos\alpha\cos\Omega.$$

**[Relationship between  $\Delta v_0$  and water content]**

The water content (in wt %) of polymer electrolyte membrane is proportional to the volume of water and therefore, to the channel volume filled with water in the membrane.  $\Delta v_0$  values are proportional to surface-to-volume ratio (S/V) of hydrophilic channels due to the anisotropic  $^2\text{H}$  interactions proportional to S/V.<sup>S2,S3</sup> Here each straight segment of channels is regarded as an individual channel so that channel tortuosity is reflected in the  $\Delta v_0$  value.

If the channel shape is assumed to be a cylinder with a radius of  $r$  and a length of  $l$ , the volume of the cylinder is  $\pi r^2 l$ , and the inner surface area of the cylinder is  $2\pi r l + 2\pi r^2$ . Consequently, S/V of the cylinder becomes  $2/r + 2/l$ . For the closed cylinders, the inner surfaces for the cylinder top and bottom are orthogonally oriented to that of the cylinder body. Thus, the anisotropic interaction proportional to the surface area of a closed cylinder requires a vector sum of  $2\pi r l$  and  $2\pi r^2$  and not a scalar sum. However, it would be more realistic to regard hydrophilic channels as cylinders with open ends. Therefore, S/V for a channel in a cylinder shape can be approximated as  $2/r$ . For rectangular columns with a length of  $l$  and widths of  $r_1$  and  $r_2$ ,  $S/V = (2lr_1 + 2lr_2 + 2r_1r_2)/(lr_1r_2) = 2/r_2 + 2/r_1 + 2/l$ . If the rectangular columns have open ends,  $S/V = 2/r_1 + 2/r_2$  and  $S/V = 4/r$  for  $r = r_1 = r_2$ . Here it can be recognized that S/V is proportional to  $\kappa/r$ , and that the proportional constant,  $\kappa$ , depends on the channel shape. If the channel shape is close to a sphere, S/V is close to  $3/r$ . But spherical shapes nullify any anisotropy, and  $\Delta v_0$  becomes zero. Even if the channel shape is not spherical, the completely randomly

oriented channels can result in zero  $\Delta v_0$  under the condition that the structures  $^2\text{H}_2\text{O}$  probed are interconnected enough for  $^2\text{H}_2\text{O}$  to sense all of the structures and exchange deuterons in the NMR time scale. The  $\Delta v_0$  values at the same water content would differ according to the degree of channel alignments. Therefore,  $\kappa$  is also a function of channel alignment or ordering degree.

When the channel volume is reduced by water content change, at least one dimensional length of the channel, for example,  $r$ , can be smaller than the anisotropic interaction length at some points. Then  $\Delta v_0$  changes to a larger value than  $\kappa$  indicates. This can be expressed such that  $\kappa$  can also be a function of the smallest dimensional length of a channel relative to the anisotropic interaction length. Here, the diffusion path lengths can be taken as an example of anisotropic interaction lengths.

The plot of  $\Delta v_0$  versus water content in wt% can be replaced with the plot of S/V of hydrophilic channels ( $\kappa/r$ ) versus the hydrophilic channel volume ( $V = n\pi r^2 l$ , where  $n$  is the number of channels) if  $\Delta v_0 = \kappa'/r$  and  $\kappa'$  is a proportional constant including the conversion constant for S/V to  $\Delta v_0$ . For simplicity, the same shape and size for all the channels were assumed. Increased  $r$  of a channel induces increased length of a neighboring channel not parallel to the channel. In addition, if the channel shape is not changed during swelling or shrinking, the ratio of  $r$  and  $l$  should be constant. Therefore,  $l$  can be replaced with  $\varepsilon r$ , where  $\varepsilon$  is a proportional constant. This results in  $V = n\pi r^2 l = n\pi \varepsilon r^3$ . The expression can be rearranged to  $r = (V/n\pi \varepsilon)^{1/3}$  and then  $\Delta v_0 = \kappa'/r$  becomes  $\Delta v_0 = \kappa'/(V/n\pi \varepsilon)^{1/3} = \kappa'(n\pi \varepsilon)^{1/3} V^{(-1/3)}$ . Thus, the plot of  $\Delta v_0$  versus water content is in the form of  $Y = \delta X^{(-1/3)}$  where  $\delta$  is a proportional constant.  $\delta = \delta_1 \delta_2 \delta_3 \delta_4$ , where  $\delta_1$  is a proportional constant related with channel alignments;  $\delta_2$  is a proportional constant related with channel shapes;  $\delta_3$  is a proportional constant related with the smallest dimensional channel length relative to the anisotropic interaction length; and  $\delta_4$  is a proportional constant related with the ratio of  $l$  to  $r$ . Among  $\delta_n$  ( $n = 1, 2, 3$ , and  $4$ ), only  $\delta_3$  can be a function of water content. When  $\delta$  is not a function of water content, the plots of  $\Delta v_0$  versus water content for different  $\delta$  values would result in the same curve shapes with different slopes.

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