

## Spectroscopic Properties and Ligand Field Analysis of $\text{trans-Dibromo(2,2-dimethyl-1,3-diaminopropane)chromium(III)}$ Moiety

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The sharp-line absorption and emission spectra of  $(\text{H}_{13}\text{O}_6)\{\text{trans-[Cr(Me}_2\text{tn)}_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)$  ( $\text{Me}_2\text{tn}$  = 2,2-dimethyl-1,3-diaminopropane) have been measured between  $13000\text{ cm}^{-1}$  and  $16000\text{ cm}^{-1}$  at 5 K. The 298 K infrared and visible absorption spectra have also been measured. The nine electronic bands due to spin-allowed and spin-forbidden transitions were assigned. Using observed transitions, a ligand field analysis has been performed to probe the ligand field properties of coordinated atoms in the title chromium(III) complex ion. The zero-phonon line in the sharp-line absorption spectrum splits into two components by  $286\text{ cm}^{-1}$ , and the large  $^2E_g$  splitting can be reproduced by the modern ligand field theory. It is confirmed that nitrogen atoms of the  $\text{Me}_2\text{tn}$  ligand have a strong  $\sigma$ -donor character, but the bromide has weak  $\sigma$ - and  $\pi$ -donor properties toward chromium(III) ion.

**Key Words :** Electronic transitions, Dibromochromium(III), Bidentate  $\text{Me}_2\text{tn}$ , AOM parameter

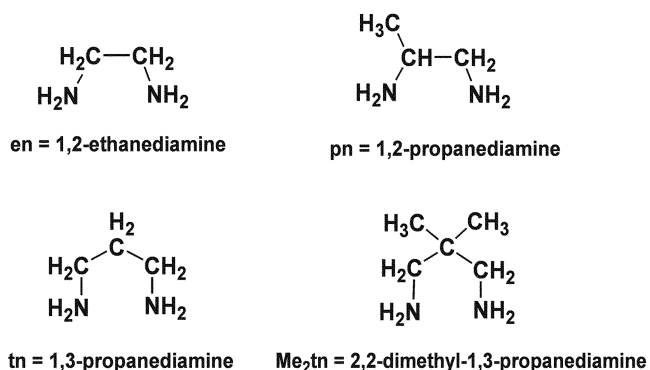
### Introduction

The chromium(III) doped crystals are promising materials for tunable solid state laser in the near-IR spectral region.<sup>1</sup> As a prerequisite for these applications, a detailed study of the spectroscopic and ligand field properties is needed. In the past two decades a considerable amount of data concerning the ligand field analysis and the highly resolved optical spectra of chromium(III) complexes have been accumulated.<sup>2-8</sup> The sharp-line electronic spectroscopic techniques offer a promise in determining metal-ligand bonding properties and in obtaining geometric information, especially in noncrystalline environments.<sup>9,10</sup>

The synthesis and structural characterization of the  $(\text{H}_{13}\text{O}_6)\{\text{trans-[Cr(Me}_2\text{tn)}_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)$  (Scheme 1;  $\text{Me}_2\text{tn}$  = 2,2-dimethyl-1,3-diaminopropane) have been reported.<sup>11</sup>

However, the vibrational and electronic energy levels based on the emission and sharp-line electronic spectroscopy of this complex have not been published yet. The structures of the  $\text{Me}_2\text{tn}$  and its related bidentate ligands are shown in Figure 1.

In this study, we have measured the low-temperature absorption and emission spectra, and the room-temperature



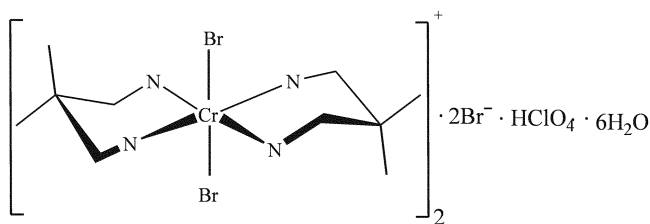
**Figure 1.** The structures of the  $\text{Me}_2\text{tn}$  and its related bidentate ligands.

infrared spectra of  $(\text{H}_{13}\text{O}_6)\{\text{trans-[Cr(Me}_2\text{tn)}_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)$ . The pure electronic origins due to spin-allowed and spin-forbidden transitions were assigned by analyzing the visible and sharp-line absorption spectra. Using the observed electronic transitions, a ligand field analysis has been performed to determine the detailed bonding properties for the coordinated bromide and nitrogen atoms toward chromium(III) ion.

### Experimental Section

**Preparations.** The ligand, 2,2-dimethyl-1,3-propanediamine was purchased from Aldrich Chemical Company. All other chemicals were of reagent grade or better quality and used without further purification. The complex  $(\text{H}_{13}\text{O}_6)\{\text{trans-[Cr(Me}_2\text{tn)}_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)$  was prepared as described previously.<sup>11</sup> The compound was recrystallized three times for spectroscopic measurements.

**Physical Measurements.** The mid-infrared spectrum was obtained with a Mattson Infinities series FT-IR spectrometer



**Scheme 1**

using a KBr pellet. The far-infrared spectrum in the region 600–50  $\text{cm}^{-1}$  was recorded with a Bruker 113v spectrometer on a microcrystalline sample pressed into a polyethylene pellet.<sup>4,5</sup> The UV-visible absorption spectrum was recorded with HP 8453 diode array spectrophotometer. The sharp-line absorption spectrum was measured on a Microcrystal spectrometer.<sup>12</sup> Luminescence spectrum was excited by using a Spectra Physics SP375 dye laser pumped by a Spectra Physics SP171  $\text{Ar}^+$  laser. The luminescence was dispersed by a Spex 1404 double monochromator and detected by a RCA 31034 photomultiplier. The samples were cooled to 5 K by a helium flow-tube cryostat.

## Results and Discussion

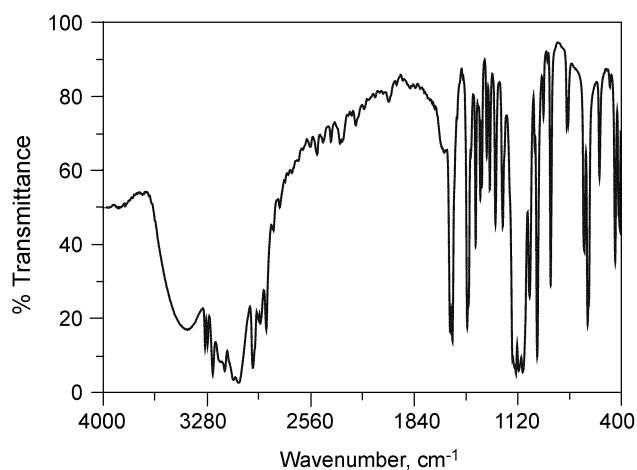
**Infrared and Emission Spectra.** The  $[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]^+$  moiety can exist as *trans* and *cis* geometric isomers. The infrared spectroscopy is frequently useful in elucidating structures and determining the number of functional groups involved in coordination. Two general approaches to the interpretation of the spectra are employed. The multiplicity of peaks associated with a given vibrational band has been used, with the less symmetric *cis* isomer having the greater multiplicity for a structurally sensitive band. The relative position of a peak in a certain absorption band has shown to be dependent on the structure of the complex.<sup>13</sup> The mid and far-infrared spectra of  $(\text{H}_{13}\text{O}_6)\{\text{trans}-[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2\cdot(\text{ClO}_4)$  recorded at room temperature are presented in Figures 2 and 3, respectively.

An absorption peak at 3422  $\text{cm}^{-1}$  can easily be assigned to the O–H stretching of  $\text{H}_2\text{O}$  molecule in the hydrate complex. The strong bands in the regions of 3100–3300  $\text{cm}^{-1}$  and 2800–3000  $\text{cm}^{-1}$  due to the symmetric and antisymmetric N–H and C–H stretching modes, respectively. The two strong absorption bands at 1590 and 1576  $\text{cm}^{-1}$  and the band at 1473  $\text{cm}^{-1}$  can be assigned to  $\text{NH}_2$  and  $\text{CH}_2$  bending modes.<sup>13</sup> The splitting of the absorption band in the 1500–1600  $\text{cm}^{-1}$  region (asymmetric  $\text{NH}_2$  deformation) may be due to the site effect. It is well known that *cis* isomer exhibits

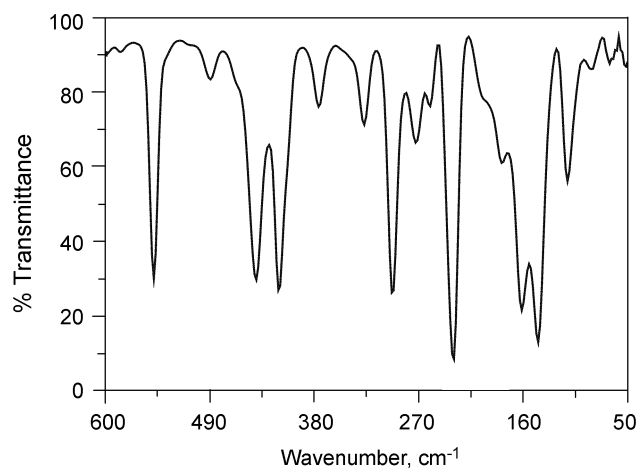
at least three bands in the 890–830  $\text{cm}^{-1}$  region due to the  $\text{NH}_2$  rocking mode while the methylene vibration splits into two peaks in the 830–760  $\text{cm}^{-1}$  region. However, *trans* isomer shows two groups of bands, one band near 890  $\text{cm}^{-1}$  arising from the amine vibration and a doublet near 800  $\text{cm}^{-1}$  due mainly to the methylene vibration.<sup>14</sup> The present complex exhibits one band at 893  $\text{cm}^{-1}$  in the  $\text{NH}_2$  rocking frequency region. The  $\text{CH}_2$  rocking bands at 771 and 780  $\text{cm}^{-1}$  are also observed. The very strong absorption at 630  $\text{cm}^{-1}$  and strong band at 1087  $\text{cm}^{-1}$  are assigned to ionic perchlorate.<sup>13</sup>

The meta-ligand stretching and ring-deformation bands occur in the far infrared range. The far-infrared spectrum (530–50  $\text{cm}^{-1}$ ) was measured in order to assign the Cr–Br and Cr–N( $\text{Me}_2\text{tn}$ ) stretching vibrations. For the *trans*- $[\text{CrN}_4\text{Br}_2]^+$  complex with  $D_{4h}$  symmetry, the group theoretical analysis predicts two Cr–N and one Cr–Br infrared-active stretching vibrations, respectively.<sup>15</sup> The two strong peaks at 442 and 419  $\text{cm}^{-1}$  can be assigned to the Cr–N stretching modes. Typically, all dibromochromium(III) complexes exhibit absorption bands in the 305–285  $\text{cm}^{-1}$  region due to the  $\nu(\text{M}-\text{Br})$ .<sup>13,15,16</sup> The Cr–Br stretching band which we have assigned a *trans* dibromo geometric configuration exhibit no splitting, whereas the *cis* isomer shows splitting of these bands, as would be expected from the lower symmetry. The one strong absorption near 299  $\text{cm}^{-1}$  may be due to the Cr–Br stretching mode. The *cis* isomer with an idealized  $C_{2v}$  symmetry exhibits two bands at 303 and 289  $\text{cm}^{-1}$ .<sup>16</sup> The title complex was consistent with the *trans* configuration. The several absorption bands lower than 274  $\text{cm}^{-1}$  arise from the skeletal bending and lattice vibration modes.

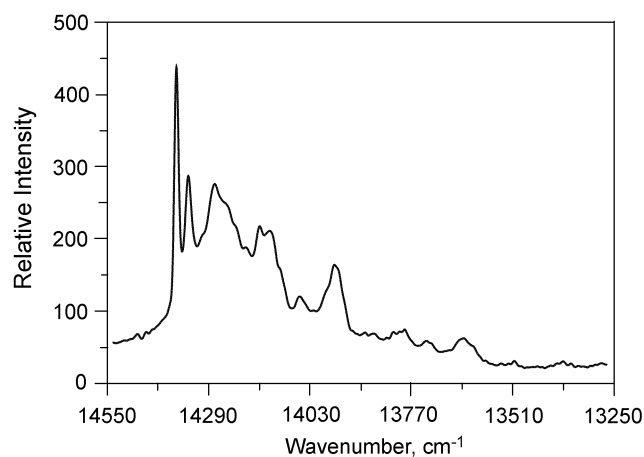
An experimental problem lies with the difficulty in distinguishing pure electronic components from the vibronic bands that also appear in the sharp-line spectrum. The emission spectra can be used in addition to the far-infrared spectra to extract the vibrational intervals of electronic ground state. The 5 K emission spectrum of  $(\text{H}_{13}\text{O}_6)\{\text{trans}-[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2\cdot(\text{ClO}_4)$  is shown in Figure 4. The band positions relative to the lowest zero phonon line, with



**Figure 2.** Mid-infrared spectrum of  $(\text{H}_{13}\text{O}_6)\{\text{trans}-[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2\cdot(\text{ClO}_4)$  at 298 K.



**Figure 3.** Far-infrared spectrum of  $(\text{H}_{13}\text{O}_6)\{\text{trans}-[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2\cdot(\text{ClO}_4)$  at 298 K.

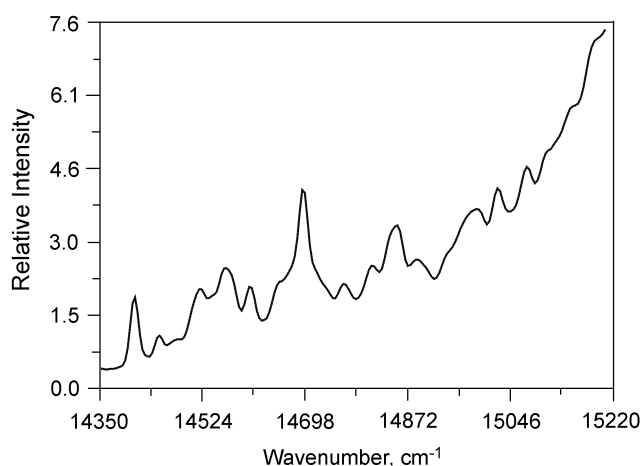


**Figure 4.** The emission spectrum of (H<sub>13</sub>O<sub>6</sub>){*trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]}<sub>2</sub>Br<sub>2</sub>(ClO<sub>4</sub>) at 5 K.

corresponding infrared frequencies, are listed in Table 1.

The very strong peak at 14374 cm<sup>-1</sup> can be assigned to the zero-phonon line, *R*<sub>1</sub> because a corresponding strong peak is found at 14413 cm<sup>-1</sup> in the absorption spectrum. The strong peak at 14343 cm<sup>-1</sup> may be assigned to the same component (*R*<sub>1</sub>') of the different site in the complex. It is conformed that the 31 cm<sup>-1</sup> splitting is due to the two inequivalent sites because the structural analysis also reveals the crystal to contain two crystallographically independent chromium(III) moieties. The bands at 419 and 477 cm<sup>-1</sup> can be assigned to a Cr-N stretching modes.<sup>17-19</sup> The vibronic intervals occurring in the spectrum consist of several modes that can be presumed to involve primarily ring torsion and angle-bending modes with frequencies in the range 250-60 cm<sup>-1</sup>.

**Absorption Spectra.** Intraconfigurational *t*<sub>2g</sub><sup>3</sup> → *t*<sub>2g</sub><sup>3</sup> transitions in chromium(III) complexes, when measured at sufficiently low temperatures, may give rise to narrow zero-



**Figure 5.** The sharp-line electronic absorption spectrum of (H<sub>13</sub>O<sub>6</sub>){*trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]}<sub>2</sub>Br<sub>2</sub>(ClO<sub>4</sub>) at 5 K.

phonon lines as well as a series of vibronic induced sharp-lines dues spin-forbidden transitions from the <sup>4</sup>*A*<sub>2g</sub> ground state into the higher doublets <sup>2</sup>*E*<sub>g</sub>, <sup>2</sup>*T*<sub>1g</sub> and <sup>2</sup>*T*<sub>2g</sub>. The 5 K sharp-line absorption spectrum of (H<sub>13</sub>O<sub>6</sub>){*trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]}<sub>2</sub>Br<sub>2</sub>(ClO<sub>4</sub>) is shown in Figure 5. The peak positions and their assignments are tabulated in Table 2. The calculated frequencies in parentheses were obtained by using the vibrational modes *v*<sub>1</sub>-*v*<sub>4</sub> listed in Table 2.

The two strong peaks at 14413 and 14699 cm<sup>-1</sup> in the sharp-line absorption spectrum are assigned to the two components (*R*<sub>1</sub> and *R*<sub>2</sub>) of the <sup>4</sup>*A*<sub>2g</sub> → <sup>2</sup>*E*<sub>g</sub> transition. The 286 cm<sup>-1</sup> splitting of the <sup>2</sup>*E*<sub>g</sub> state can be compared to the 198 cm<sup>-1</sup> and 181 cm<sup>-1</sup> observed for *cis*-[Cr(trien)Cl<sub>2</sub>]Cl and [Cr(NH<sub>3</sub>)<sub>5</sub>(imH)](ClO<sub>4</sub>)<sub>3</sub> complexes, respectively.<sup>20,21</sup> The medium peak at 14457 cm<sup>-1</sup> may be assigned to the same component (*R*<sub>1</sub>') of the different site in the complex thus the 44 cm<sup>-1</sup> splitting is due to the site effect.

In general, it is not easy to locate the positions of the other electronic components because the vibronic sidebands of the <sup>2</sup>*E*<sub>g</sub> levels overlap with the zero phonon lines of <sup>2</sup>*T*<sub>1g</sub>. However, the three tentative components of the <sup>4</sup>*A*<sub>2g</sub> → <sup>2</sup>*T*<sub>1g</sub> electronic origin (*T*<sub>1</sub>, *T*<sub>2</sub> and *T*<sub>3</sub>) are assigned to relatively intense peaks at 580, 618 and 660 cm<sup>-1</sup> from the lowest electronic line, *R*<sub>1</sub>. Vibronic satellites based on these origins have similar frequencies and intensity patterns to those of the <sup>2</sup>*E*<sub>g</sub> components. These three electronic lines are somewhat less intense than the <sup>2</sup>*E*<sub>g</sub> lines, but they are relatively intense compared to other chromium(III) complexes.

The visible absorption spectrum (solid line) of (H<sub>13</sub>O<sub>6</sub>)-{*trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]}<sub>2</sub>Br<sub>2</sub>(ClO<sub>4</sub>) in aqueous solution at room temperature is also represented in Figure 6.

It exhibits main two bands, one at 16180 cm<sup>-1</sup> (*v*<sub>1</sub>), and the other at 25810 cm<sup>-1</sup> (*v*<sub>2</sub>), corresponding to the <sup>4</sup>*A*<sub>2g</sub> → <sup>4</sup>*T*<sub>2g</sub> and <sup>4</sup>*A*<sub>2g</sub> → <sup>4</sup>*T*<sub>1g</sub> (*O*<sub>h</sub>) transitions, respectively.<sup>22</sup> The assignment of geometrical configuration is suggested by inspection of the d-d absorption spectra. The position of the spin-allowed transitions in the electronic spectra, the number of bands, and their coefficients are usually reliable indicators for distinguishing the *cis* and *trans* geometrical isomers. In

**Table 1.** Vibrational frequencies from the 5 K Emission and 298 K infrared spectra for (H<sub>13</sub>O<sub>6</sub>){*trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]}<sub>2</sub>Br<sub>2</sub>(ClO<sub>4</sub>)<sup>a</sup>

Emission <sup>b</sup>	Infrared	Assignment
0 vs		<i>R</i> <sub>1</sub>
31 s		<i>R</i> <sub>1</sub> '
98 m	68 w, 89 sh	Lattice vib., skeletal bends and <i>v</i> (Cr-Br)
133 sh	114 s, 145 m	
214 m	183 vw, 214 sh	
244 m	233 vs, 258 vw, 274 w 299 vs	
317 w	328 m, 376 w	<i>v</i> (Cr-N)
419 s	419 vs, 442 vs	
477 vw	490 vw	
506 vw		<i>v</i> (Cr-N) + ring def.
562 w	550 vs	
645 w	630 vs, 659 s	ClO <sub>4</sub> <sup>-</sup>
743 w	770 m, 780 m	<i>ρ</i> (CH <sub>2</sub> )
866 w		
891 w	893 vs	<i>ρ</i> (NH <sub>2</sub> )

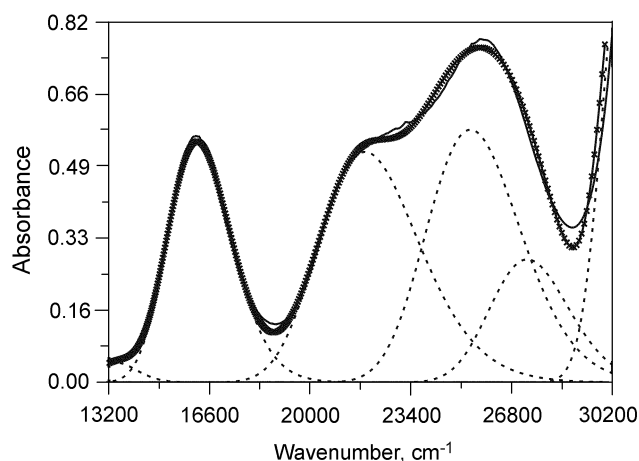
<sup>a</sup>Data in cm<sup>-1</sup>. <sup>b</sup>Measured from zero-phonon line at 14374 cm<sup>-1</sup>.

**Table 2.** Peak positions in the 5 K sharp-line absorption spectrum of  $(\text{H}_{13}\text{O}_6)\{\text{trans}[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)^a$ 

$\bar{\nu}_0 - 14413$	Assignment	Calcd <sup>b</sup>	Vibronic frequencies	Ground state frequencies <sup>c</sup>
0 vs	$R_1$		$\nu_1$ 117	114
44 m	$R_1'$		$\nu_2$ 157	162
111 m	$R_1 + \nu_1$	(117)	$\nu_3$ 252	258
160 s	$R_1 + \nu_2$	(157)	$\nu_4$ 444	444
196 s	$R_1 + \nu_2$	(201)		
258 w	$R_1 + \nu_3$	(252)		
286 vs	$R_2$			
321 sh	$R_1 + 2\nu_1$	(314)		
358 m	$R_1' + 2\nu_1$	(358)		
406 m	$R_2 + \nu_1$	(403)		
444 vs	$R_1 + \nu_4$	(444)		
479 m	$R_2 + \nu_4$	(488)		
580 s	$T_1$			
618 vs	$T_2$			
668 s	$T_3$			
703 w	$T_1 + \nu_1$	(697)		
739 w	$T_1 + \nu_2$	(737)		
783 w	$T_3 + \nu_3$	(785)		
833 sh	$T_1 + \nu_2$	(832)		

<sup>a</sup>Data in  $\text{cm}^{-1}$ . <sup>b</sup>Values in parentheses represent the calculated frequencies based on the vibrational modes listed. <sup>c</sup>From the emission and far-infrared spectra (Table 1).

general, the less symmetrical *cis* chromium(III) complexes have two symmetric bands in the visible region, and these bands are located at lower wavelengths and have higher extinction coefficients than those of more symmetrical *trans*-isomers.<sup>5-7,22</sup> The one quartet band has a shoulder peak at  $21740\text{ cm}^{-1}$  and shows asymmetric profile. This property of the quartet band is indicative of the *trans* configuration. In order to have some point of the reference for the splitting of the bands, we have fit the band profiles to four Gaussian curves, as seen in Figure 5. The contribution from outside bands was corrected for fine deconvolution.<sup>23</sup> A deconvolution procedure<sup>21</sup> on the experimental band pattern yielded

**Figure 6.** The visible absorption spectrum of  $\text{trans}[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]^+$  and the resolved overlapping peaks (dotted curves).

maxima at  $16170$ ,  $21845$ ,  $25400$  and  $27320\text{ cm}^{-1}$  for the noncubic split levels of  ${}^4T_{2g}$  and  ${}^4T_{1g}$ , respectively. These resolved peak positions were used as the observed spin-allowed transition energies in the ligand field optimization. In fact, using just one Gaussian curve instead of two yields a least squares error only four times that of the best fit (dotted line) shown in Figure 6.

**AOM Calculations.** All ligand field calculations are carried out using the AOMX program.<sup>24</sup> A full ligand field calculation on chromium(III) complexes requires the diagonalization of a 120-dimensional secular determinant which arises from the perturbed  $d^3$  system. The ligand field potential matrix was generated for  $\text{trans}[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]^+$  moiety from the coordinated four nitrogen and two bromide atoms. In the framework of the angular overlap model the metal-ligand interactions are described in terms of localized bonding parameters of  $\sigma$ - and  $\pi$ -type. The parameters varied during the optimization were the interelectronic repulsion parameters  $B$ ,  $C$  and the Trees correction parameter  $\tau$ , the spin-orbit coupling parameter  $\zeta$ , the AOM parameters  $e_\sigma(\text{Br})$  and  $e_\pi(\text{Br})$  for the bromide-chromium, and  $e_\sigma(\text{N})$  for the  $\text{Me}_2\text{tn}$  nitrogen-chromium. The  $\pi$ -interaction of the amine nitrogen with  $sp^3$  hybridization in the  $\text{Me}_2\text{tn}$  was assumed to be negligible. However, it is noteworthy that the peptide nitrogen with  $sp^2$  hybridization has a weak  $\pi$ -donor character.<sup>25-28</sup> Schmidtke's  $\pi$ -expansion parameter  $\tau$  were also included in the treatment of the interelectronic repulsion term.<sup>3</sup> In Schmidtke's approximation, the electrostatic terms are modified by a factor  $\tau$  for each constituent metal wavefunction that overlaps with a ligand  $\pi$ -orbital. The  $\pi$ -orbital expansion parameter,  $\tau$  was optimized at the value 0.9892. Racah parameters depend on the value of Trees parameter  $a_T$ , which was included in the calculations.<sup>29</sup> For calculating minor doublet splitting, the effect of the spin-orbit coupling may become important and has to be considered then by introducing an additional parameter  $\zeta$ . The angular positions of the ligating six atoms were taken from the X-ray crystal structure<sup>11</sup> of  $(\text{H}_{13}\text{O}_6)\{\text{trans}[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]\}_2\text{Br}_2(\text{ClO}_4)$ , which was determined to be the space group  $C2/m$  of the monoclinic system with a cell of dimensions  $a = 22.3862(8)$ ,  $b = 24.7393(8)$ ,  $c = 8.3230(3)\text{ \AA}$  and  $\beta = 91.6240(10)^\circ$ . The coordinates were then rotated so as to maximize the projections of the six-coordinated atoms on the Cartesian axes centered on the chromium. The resulting Cartesian and spherical coordinates for the  $\text{trans}[\text{Cr}(\text{Me}_2\text{tn})_2\text{Br}_2]^+$  moiety are listed in Table 3.

The ligand field analysis was carried out through an optimized fit of experimental to calculated transition energies as follows. Since the Racah parameters influence the low-symmetry level splitting only by high order perturbation, these quantities can be determined in good approximation from the quartet band separation ( ${}^4T_{1g} - {}^4T_{2g} \approx 12B$ ) and from the energy position of lowest doublet term ( ${}^2E_g - {}^4A_{2g} \approx 9B + 3C$ ). Moreover, AOM parameters should be transferable between similar compounds, leading to a value close to  $7300\text{ cm}^{-1}$  for  $e_\sigma(\text{N})$  parameter.<sup>2-8</sup> All parameters, except  $e_\sigma(\text{Br})$  and  $e_\pi(\text{Br})$ , were constrained to reasonable

**Table 3.** Optimized Cartesian and spherical polar coordinates for ligating atoms in *trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]<sup>+</sup> moiety<sup>a</sup>

Atom <sup>b</sup>	x	y	z	θ	φ
N <sub>11</sub>	-0.0267	-2.0781	0.0108	89.70	-90.74
N <sub>12</sub>	2.0824	0.0285	0.0042	89.88	0.78
N <sub>11</sub> <sup>i</sup>	0.0267	2.0781	-0.0109	90.30	89.26
N <sub>12</sub> <sup>i</sup>	-2.0868	-0.0248	-0.0078	90.21	-179.32
Br <sub>11</sub>	0.0050	-0.0091	2.4878	0.24	-61.21
Br <sub>11</sub> <sup>i</sup>	-0.0050	0.0090	-2.4879	179.76	119.06

<sup>a</sup>Cartesian coordinates in Å, polar coordinates in degrees. <sup>b</sup>Atomic labeling was adopted from Ref. 11.

limits based on the data from the other chromium(III) complexes. The six parameters were used to fit ten experimental energies: the five  $^4A_{2g} \rightarrow \{^2E_g, ^2T_{1g}\}$  components, identified in Table 4, the four  $^4A_{2g} \rightarrow \{^4T_{2g}, ^4T_{1g}\}$  components, and the splitting of the  $^2E_g$  state. The Powell parallel subspace optimization procedure<sup>30</sup> was used to find the global minimum. The optimization was repeated several times with the different sets of starting parameters to verify that the same global minimum was found. The results of the optimization and the parameter set used to generate the best-fit energies are also listed in Table 4. The error margins reported for the best-fit parameters in Table 4 are based only on the propagation of the assumed uncertainties in the observed peak positions.<sup>31</sup> The quartet terms were given a very low weight to reflect the very large uncertainty in their position.

The deduced ligand field parameters are  $e_\sigma(\text{N}) = 7404 \pm 23$ ;  $e_\sigma(\text{Br}) = 4879 \pm 48$ ;  $e_\pi(\text{Br}) = 797 \pm 16$ ;  $B = 771 \pm 9$ ;  $C = 3115 \pm 29$ ;  $\alpha_T = 200 \pm 10$  and  $z = 149 \pm 45 \text{ cm}^{-1}$ . The AOM parameters are plausible and reproduce the spectrum pretty well. The  $e_\sigma(\text{N})$  value for amine nitrogen is located in the normal range.<sup>32-37</sup> A ligand field analysis of the sharp-line absorption and broad-band absorption spectra indicates that the bromide is a weak  $\sigma$ - and  $\pi$ -donor. The value of  $7404 \text{ cm}^{-1}$  for  $e_\sigma(\text{N})$  is comparable to values for the other amines.<sup>17-19</sup> It is suggested that the four nitrogen atoms of the bidentate Me<sub>2</sub>tn ligand have strong  $\sigma$ -donor properties

**Table 4.** Experimental and calculated electronic transition energies for *trans*-[Cr(Me<sub>2</sub>tn)<sub>2</sub>Br<sub>2</sub>]<sup>+</sup> moiety<sup>a</sup>

State ( <i>O<sub>h</sub></i> )	Exptl	Calcd <sup>b</sup>
$^2E_g$	14413	14407
	14699	14696
$^2T_{1g}$	14993	14937
	15031	15109
$^4T_{2g}$	15081	15179
	16170 <sup>c</sup>	16221
$^4T_{1g}$	21845 <sup>c</sup>	21966
	25400 <sup>c</sup>	24697
	27320 <sup>c</sup>	27202

<sup>a</sup>Data in  $\text{cm}^{-1}$ . <sup>b</sup> $e_\sigma(\text{N}) = 7404 \pm 23$ ;  $e_\sigma(\text{Br}) = 4879 \pm 48$ ;  $e_\pi(\text{Br}) = 797 \pm 16$ ;  $B = 771 \pm 9$ ;  $C = 3115 \pm 29$ ;  $\alpha_T = 200 \pm 10$ ;  $z = 149 \pm 45$ . <sup>c</sup>Obtained from the Gaussian component deconvolution.

toward chromium(III). An orbital population analysis yields a configuration of  $(xy)^{1.000}(xz)^{0.972}(yz)^{0.942}(x^2-y^2)^{0.017}(z^2)^{0.069}$  for the lowest component of the  $^2E_g$  state. The relative *d*-orbital ordering from the calculation is  $E(xy) = 8.3 \text{ cm}^{-1} < E(xz) = 1593 \text{ cm}^{-1} < E(yz) = 1595 \text{ cm}^{-1} < E(x^2-y^2) = 17167 \text{ cm}^{-1} < E(z^2) = 22198 \text{ cm}^{-1}$ . The value of Racah parameter, *B*, is about 84% of the value for a free chromium(III) ion in the gas phase. These factors plus AOM parameters can be used in determining the photolabilized ligand of the photoreaction in the chromium(III) complexes and predicting the relative efficiency of the 3d-4f energy transfer in the heterometal dinuclear complexes.<sup>38,39</sup> The  $286 \text{ cm}^{-1}$  of observed  $^2E_g$  splitting in the sharp-line absorption spectrum is slightly larger than the  $210$  and  $239 \text{ cm}^{-1}$  of *trans*-[Cr(en)<sub>2</sub>Cl<sub>2</sub>]Cl and *trans*-[Cr(tn)<sub>2</sub>Cl<sub>2</sub>]ClO<sub>4</sub>, respectively.<sup>40</sup> It is shown that the large  $^2E_g$  splitting can be reproduced by the modern ligand field theory. The parameter values reported here appear to be significant, as deduced on the basis of the manifold of sharp line transitions in which were obtained from the well resolved absorption spectra. The infrared and electronic spectral properties of this complex are in good agreement with the result<sup>11</sup> of X-ray crystallography, which show that the Cr atom is an octahedral environment coordinated by two bidentate 2,2-dimethyl-1,3-diaminopropanes and two Br atoms in a *trans* positions.

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